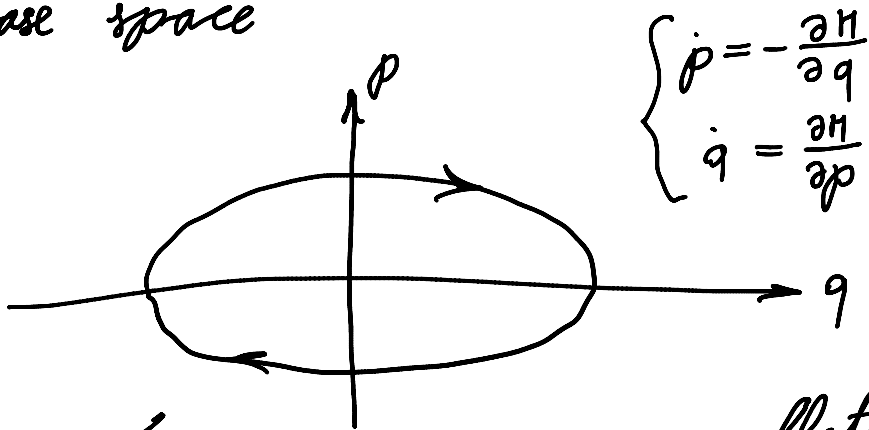


## Concept of thermalisation. Liouville's theorem

We've learnt of the concept of phase space. For a closed system, motion occurs along deterministic trajectories in phase space



Phase trajectory of an oscillator

$$H = \frac{p^2}{2m} + \frac{m\omega^2 q^2}{2}$$

If the system is connected to an environment, it may "forget" its state after some time  $T_{\text{relax}}$ .

Equilibrium = the state of the system is independent of how the system came there

Environment = thermostat = bath = the environment which brings the system to equilibrium

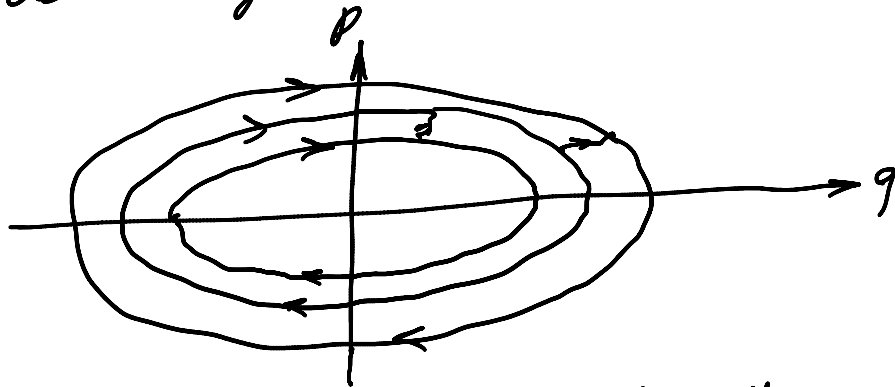
Example

## Example

$$H = \underbrace{\frac{p^2}{2m} + \frac{m\omega^2 q^2}{2}}_{H_{\text{osc.}}} + \underbrace{g\xi}_{V_{\text{coupling}}} + \underbrace{H_b(\xi, P_\xi)}_{\text{Hamiltonian of the bath}}$$

For a weak coupling, the system will come to equilibrium on a long time  $T_{\text{relax}}$ .

At times  $t \ll T_{\text{relax}}$  the system still obeys the equations of motions of a closed system



Coupling to the environment makes the system ergodic = makes the trajectories dense in phase space

Let's find how the density evolves

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## Liouville's theorem

Instead of time average, one may think about ensemble averages. Then  $\rho$  is the density of systems

ensemble averages. Then  $\rho$  is the density of  $\Gamma$  in phase space.

$$\frac{\partial \rho}{\partial t} + \text{div}(\rho \vec{v}) = 0 \quad (\text{For isolated systems})$$

$$\text{here } \text{div}(\rho \vec{v}) = \sum_{i=1}^f \left( \frac{\partial}{\partial q_i} (\rho \dot{q}_i) + \frac{\partial}{\partial p_i} (\rho \dot{p}_i) \right)$$

$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho \dot{q})}{\partial \dot{q}} + \frac{\partial (\rho \dot{p})}{\partial \dot{p}} = 0$$

$$\rho \frac{\partial \dot{q}}{\partial \dot{q}} + \rho \frac{\partial \dot{p}}{\partial \dot{p}} = \rho \frac{\partial^2 H}{\partial \dot{q} \partial \dot{p}} - \rho \frac{\partial^2 H}{\partial \dot{q} \partial \dot{p}} = 0$$

$$\rightarrow \frac{d\rho}{dt} = \frac{\partial \rho}{\partial t} + \frac{\partial \rho}{\partial \dot{q}} \dot{q} + \frac{\partial \rho}{\partial \dot{p}} \dot{p} = 0$$

In an isolated system the distribution function is constant along phase trajectories